

QUANTUM INTEGRABILITY TOWARDS SUPERCONDUCTIVITY

Casey Farren-Collopy

Supervised by: Sergey Frolov

Institute: TCD School of Maths

GIVEN TO THE SCHOOL OF MATHS STUDENTS AS PART OF THE CAPSTONE PRESENTATIONS 2025/26.

Table of contents

Quantum Integrability

Bethe Ansatz

The Hubbard Model

References

Appendices

Quantum Integrability

- © **Quantum Groups** underpin the mathematics of integrability - but they're awful[2]

- ⊙ **Quantum Groups** underpin the mathematics of integrability - but they're awful[2]
- ⊙ They are defined around a central object - the R matrix

Quantum Groups

- ⊙ **Quantum Groups** underpin the mathematics of integrability - but they're awful[2]
- ⊙ They are defined around a central object - the R matrix
- ⊙ Each system is equivalent to a representation of the universal R matrix

Quantum Groups

- ⊙ **Quantum Groups** underpin the mathematics of integrability - but they're awful[2]
- ⊙ They are defined around a central object - the R matrix
- ⊙ Each system is equivalent to a representation of the universal R matrix
- ⊙ Physically, the R matrix describes interactions between a pair of particles

Quantum Groups

- ⊙ **Quantum Groups** underpin the mathematics of integrability - but they're awful[2]
- ⊙ They are defined around a central object - the R matrix
- ⊙ Each system is equivalent to a representation of the universal R matrix
- ⊙ Physically, the R matrix describes interactions between a pair of particles
- ⊙ There are two relations defining a **Quantum Group**

Quantum Groups

- ⊙ **Quantum Groups** underpin the mathematics of integrability - but they're awful[2]
- ⊙ They are defined around a central object - the R matrix
- ⊙ Each system is equivalent to a representation of the universal R matrix
- ⊙ Physically, the R matrix describes interactions between a pair of particles
- ⊙ There are two relations defining a **Quantum Group**

1. Yang-Baxter Equation

Quantum Groups

- ⊙ **Quantum Groups** underpin the mathematics of integrability - but they're awful[2]
- ⊙ They are defined around a central object - the R matrix
- ⊙ Each system is equivalent to a representation of the universal R matrix
- ⊙ Physically, the R matrix describes interactions between a pair of particles
- ⊙ There are two relations defining a **Quantum Group**
 1. Yang-Baxter Equation
 2. RTT Relations

The Basis of Quantum Integrability

Yang-Baxter

$$R_{12}(\lambda, \mu)R_{13}(\lambda, \tau)R_{23}(\mu, \tau) \\ = R_{23}(\mu, \tau)R_{13}(\lambda, \tau)R_{12}(\lambda, \mu).$$

There are two contexts where we can physically interpret it:

1. Scattering

^aOne dimensional lattices like the Heisenberg spin chain.

The Basis of Quantum Integrability

Yang-Baxter

$$R_{12}(\lambda, \mu)R_{13}(\lambda, \tau)R_{23}(\mu, \tau) \\ = R_{23}(\mu, \tau)R_{13}(\lambda, \tau)R_{12}(\lambda, \mu).$$

There are two contexts where we can physically interpret it:

1. Scattering
2. Chains^a

^aOne dimensional lattices like the Heisenberg spin chain.

The Basis of Quantum Integrability

Yang-Baxter

$$\begin{aligned} & R_{12}(\lambda, \mu)R_{13}(\lambda, \tau)R_{23}(\mu, \tau) \\ &= R_{23}(\mu, \tau)R_{13}(\lambda, \tau)R_{12}(\lambda, \mu). \end{aligned}$$

There are two contexts where we can physically interpret it:

1. Scattering
2. Chains^a

We will focus on the latter.

^aOne dimensional lattices like the Heisenberg spin chain.

The Basis of Quantum Integrability


Yang-Baxter

$$R_{12}(\lambda, \mu)R_{13}(\lambda, \tau)R_{23}(\mu, \tau) \\ = R_{23}(\mu, \tau)R_{13}(\lambda, \tau)R_{12}(\lambda, \mu).$$

There are two contexts where we can physically interpret (4):

1. Scattering
2. Chains^a

We will focus on the latter.

 **Problem:** Chains have all local interactions between particles 1 & 3 doesn't really make sense.

^aOne dimensional lattices like the Heisenberg spin chain.

The Basis of Quantum Integrability


Yang-Baxter

$$R_{12}(\lambda, \mu)R_{13}(\lambda, \tau)R_{23}(\mu, \tau) \\ = R_{23}(\mu, \tau)R_{13}(\lambda, \tau)R_{12}(\lambda, \mu).$$

There are two contexts where we can physically interpret (4):

1. Scattering
2. Chains^a

We will focus on the latter.

 **Problem:** Chains have all local interactions between particles 1 & 3 doesn't really make sense.

 **Solution:** We *twist* the R matrix. [3]

^aOne dimensional lattices like the Heisenberg spin chain.

The Basis of Quantum Integrability

Yang-Baxter


$$R_{12}(\lambda, \mu)R_{13}(\lambda, \tau)R_{23}(\mu, \tau) \\ = R_{23}(\mu, \tau)R_{13}(\lambda, \tau)R_{12}(\lambda, \mu).$$

There are two contexts where we can physically interpret (4):

1. Scattering
2. Chains^a

We will focus on the latter.

^aOne dimensional lattices like the Heisenberg spin chain.

 **Problem:** Chains have all local interactions between particles 1 & 3 doesn't really make sense.

 **Solution:** We *twist* the R matrix. [3]

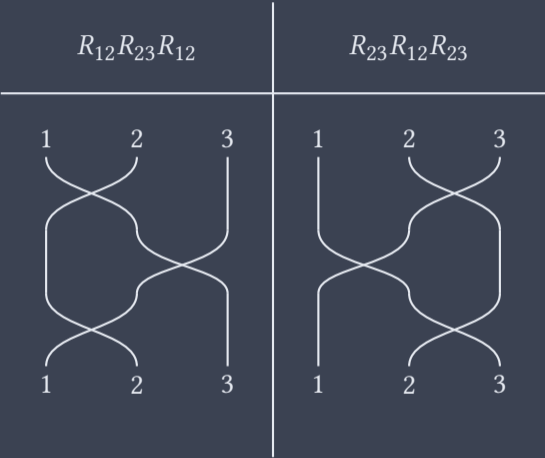
Braid Relation

Twisting the R matrix gets us a new but equivalent equation,

$$R_{12}R_{23}R_{12} = R_{23}R_{12}R_{23} \quad (1)$$

But what does it mean?

The Basis of Quantum Integrability >_ But what does it mean?



CONSISTENCY!

The T Matrix & Conserved Quantities

The RTT Relations

$$R_{ab}(\lambda, \mu)T_a(\lambda)T_b(\mu) = T_b(\mu)T_a(\lambda)R_{ab}(\lambda, \mu).$$

The T Matrix & Conserved Quantities

The RTT Relations

$$R_{ab}(\lambda, \mu)T_a(\lambda)T_b(\mu) = T_b(\mu)T_a(\lambda)R_{ab}(\lambda, \mu).$$

Describes a fictitious particle's effect on the entire system as is seen by the solution,

$$T_a(\lambda) = \prod_{k=1}^N R_{ka}(p_k, \lambda). \quad (2)$$

The T Matrix & Conserved Quantities

The RTT Relations

$$R_{ab}(\lambda, \mu)T_a(\lambda)T_b(\mu) = T_b(\mu)T_a(\lambda)R_{ab}(\lambda, \mu).$$

Describes a fictitious particle's effect on the entire system as is seen by the solution,,

$$T_a(\lambda) = \prod_{k=1}^N R_{ka}(p_k, \lambda). \quad (3)$$

⊙ Need to introduce one more thing...

The T Matrix & Conserved Quantities

The RTT Relations

$$R_{ab}(\lambda, \mu)T_a(\lambda)T_b(\mu) = T_b(\mu)T_a(\lambda)R_{ab}(\lambda, \mu).$$

Describes a fictitious particle's effect on the entire system as is seen by the solution,,

$$T_a(\lambda) = \prod_{k=1}^N R_{ka}(p_k, \lambda). \quad (3)$$

⊙ Need to introduce one more thing...

The Transfer Matrix

We define,

$$\tau(p) := \text{Tr}_a\{T_a(p)\} \quad \begin{array}{c} \text{It can be shown} \\ \implies \end{array} \quad [\tau(p), \tau(k)] = 0$$

But why is this so useful?

The Bethe Ansatz(s)

- ⊙ There's a bit of ambiguity in saying **Bethe Ansatz**

The Bethe Ansatz(s)

- ⊙ There's a bit of ambiguity in saying **Bethe Ansatz**
- ⊙ One of two techniques:

The Bethe Ansatz(s)

- ⊙ There's a bit of ambiguity in saying **Bethe Ansatz**
- ⊙ One of two techniques:
 1. Coordinate Bethe Ansatz

The Bethe Ansatz(s)

- ⊙ There's a bit of ambiguity in saying **Bethe Ansatz**
- ⊙ One of two techniques:
 1. Coordinate Bethe Ansatz
 2. Algebraic Bethe Ansatz

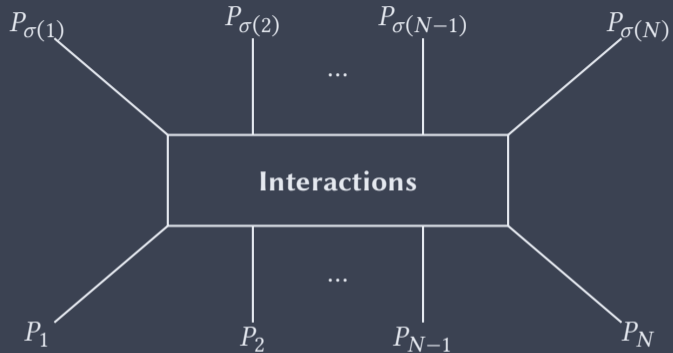
The Bethe Ansatz(s)

- ⊙ There's a bit of ambiguity in saying **Bethe Ansatz**
- ⊙ One of two techniques:
 1. Coordinate Bethe Ansatz
 2. Algebraic Bethe Ansatz

Factorised Scattering


The assumption of **factorised scattering** applied to a system is called the **Coordinate Bethe Ansatz**.

Factorised Scattering




Note

Here $\sigma \in \mathfrak{S}^N$ is some permutation since momentum conservation reflects a non-diffractive scattering process.

 **Problem:** Often the Coordinate Bethe Ansatz is quite limited - but there exists a far more powerful tool.


Algebraic Bethe Ansatz

 **Problem:** Coordinate Bethe Ansatz is quite limited - but there exists a far more powerful tool.

The Algebraic Bethe Ansatz


1. **Quantum Group** representation $\rightarrow T$ matrix.

Algebraic Bethe Ansatz

 **Problem:** Coordinate Bethe Ansatz is quite limited - but there exists a far more powerful tool.

The Algebraic Bethe Ansatz


1. **Quantum Group** representation $\rightarrow T$ matrix.
2. Find *classes* of states.

 **Problem:** Coordinate Bethe Ansatz is quite limited - but there exists a far more powerful tool.

The Algebraic Bethe Ansatz

1. **Quantum Group** representation $\rightarrow T$ matrix.
2. Find *classes* of states.
3. Raising and lowering operators built from symmetries extends state space.

Algebraic Bethe Ansatz


 **Problem:** Coordinate Bethe Ansatz is quite limited - but there exists a far more powerful tool.

The Algebraic Bethe Ansatz

1. **Quantum Group** representation $\rightarrow T$ matrix.
2. Find *classes* of states.
3. Raising and lowering operators built from symmetries extends state space.

Reminds us of Clebsch-Gordan.

Algebraic Bethe Ansatz

 **Problem:** Coordinate Bethe Ansatz is quite limited - but there exists a far more powerful tool.


The Algebraic Bethe Ansatz

1. **Quantum Group** representation $\rightarrow T$ matrix.
2. Find *classes* of states.
3. Raising and lowering operators built from symmetries extends state space.

Reminds us of Clebsch-Gordan.

What's the difference between the two?

Algebraic Bethe Ansatz

 **Problem:** Coordinate Bethe Ansatz is quite limited - but there exists a far more powerful tool.

The Algebraic Bethe Ansatz

1. **Quantum Group** representation $\rightarrow T$ matrix.
2. Find *classes* of states.
3. Raising and lowering operators built from symmetries extends state space.

Reminds us of Clebsch-Gordan.

What's the difference between the two?

E.g. Hydrogen Atom.[4]

The Hubbard Model

Spectrum

- ⊙ Applying the coordinate Bethe ansatz one sees that the charge and spin degrees of freedom separate
- ⊙ The charge degrees of freedom called **holons** and spin degrees are called **spinons**.

Similarly we get the spectrum,

$$E = -2 \sum_j \cos(k_j) + u(L - 2N), \quad P = \sum_j k_j \pmod{2\pi}$$

Spectrum

- ⊙ Applying the coordinate Bethe ansatz one sees that the charge and spin degrees of freedom separate
- ⊙ The charge degrees of freedom called **holons** and spin degrees are called **spinons**.

Similarly we get the spectrum,

$$E = -2 \sum_j \cos(k_j) + u(L - 2N), \quad P = \sum_j k_j \pmod{2\pi}$$

Note here what happens with half-filling.

So what about superconductivity?

- ⊙ We apply the Algebraic Bethe Ansatz

So what about superconductivity?

- ⊙ We apply the Algebraic Bethe Ansatz
- ⊙ This manifests an $SO(4) \simeq \frac{SU(2) \times SU(2)}{\mathbb{Z}_2}$ symmetry

So what about superconductivity?

- ⊙ We apply the Algebraic Bethe Ansatz
- ⊙ This manifests an $SO(4) \simeq \frac{SU(2) \times SU(2)}{\mathbb{Z}_2}$ symmetry
- ⊙ One $SU(2)$ corresponds to the spin DoF and the other the charge - called η pairs

So what about superconductivity?

- ⊙ We apply the Algebraic Bethe Ansatz
- ⊙ This manifests an $SO(4) \simeq \frac{SU(2) \times SU(2)}{\mathbb{Z}_2}$ symmetry
- ⊙ One $SU(2)$ corresponds to the spin DoF and the other the charge - called η pairs

However...

So what about superconductivity?

- ⊙ We apply the Algebraic Bethe Ansatz
- ⊙ This manifests an $SO(4) \simeq \frac{SU(2) \times SU(2)}{\mathbb{Z}_2}$ symmetry
- ⊙ One $SU(2)$ corresponds to the spin DoF and the other the charge - called η pairs

However...

 **Problem:** The η pairs cause the decay of superconducting modes.

So what about superconductivity?

- ⊙ We apply the Algebraic Bethe Ansatz
- ⊙ This manifests an $SO(4) \simeq \frac{SU(2) \times SU(2)}{\mathbb{Z}_2}$ symmetry
- ⊙ One $SU(2)$ corresponds to the spin DoF and the other the charge - called η pairs

However...

 **Problem:** The η pairs cause the decay of superconducting modes.

 **Solution:** Go to two dimensions - and give up integrability...

Thank You For Listening!

There are a lot more interesting & relevant points in condensed matter that I did not get to cover. For instance:

Thank You For Listening!

There are a lot more interesting & relevant points in condensed matter that I did not get to cover. For instance:

- ⊙ The Quantum Hall Effect

Thank You For Listening!

There are a lot more interesting & relevant points in condensed matter that I did not get to cover. For instance:

- ⊙ The Quantum Hall Effect
- ⊙ The Hubbard Model as a Mott Insulator

Thank You For Listening!

There are a lot more interesting & relevant points in condensed matter that I did not get to cover. For instance:

- ⊙ The Quantum Hall Effect
- ⊙ The Hubbard Model as a Mott Insulator

But in any case,

Thank You For Listening!

There are a lot more interesting & relevant points in condensed matter that I did not get to cover. For instance:

- ⊙ The Quantum Hall Effect
- ⊙ The Hubbard Model as a Mott Insulator

But in any case,

Thanks for listening!

References

- [1] G. Arutyunov. *Elements of Classical and Quantum Integrable Systems*. UNITEXT for Physics. Springer International Publishing, 2019. ISBN: 9783030241971. URL: <https://books.google.ie/books?id=5sPhxQEACAAJ>.
- [2] Vyjayanthi Chari and Andrew Pressley. *A guide to quantum groups*. Reprint. Cambridge: Cambridge Univ. Press, 2000. 651 pp. ISBN: 9780521558846 9780521433051.
- [3] Larry A. Lambe and David E. Radford. *Introduction to the Quantum Yang-Baxter Equation and Quantum Groups: An Algebraic Approach*. Boston, MA: Springer US, 1997. ISBN: 9781461368427 9781461541097. DOI: 10.1007/978-1-4615-4109-7. URL: <http://link.springer.com/10.1007/978-1-4615-4109-7> (visited on 01/26/2026).
- [4] W. Pauli. “Über das Wasserstoffspektrum vom Standpunkt der neuen Quantenmechanik”. In: *Zeitschrift für Physik A Hadrons and nuclei* 36.5 (May 1926), pp. 336–363. ISSN: 0939-7922. DOI: 10.1007/BF01450175. URL: <http://link.springer.com/10.1007/BF01450175> (visited on 03/08/2026).

(Quantum) R matrix

- ⊙ We¹ can think of the R matrix as the object that describes a pair of particles with two momenta.

¹As physicists.

(Quantum) R matrix

- ⊙ We¹ can think of the R matrix as the object that describes a pair of particles with two momenta.
- ⊙ If we say that the i^{th} particle lives in \mathcal{H}_i then the system lives in $\mathcal{H}_{\text{sys}} := \bigotimes_{k=1}^N \mathcal{H}_i$

¹As physicists.


(Quantum) R matrix

- ⊙ We¹ can think of the R matrix as the object that describes a pair of particles with two momenta.
- ⊙ If we say that the i^{th} particle lives in \mathcal{H}_i then the system lives in $\mathcal{H}_{\text{sys}} := \bigotimes_{k=1}^N \mathcal{H}_i$
- ⊙ The R matrix then is defined on two copies of \mathcal{H} - the state space of the two particle it describes.

¹As physicists.

(Quantum) R matrix

- ⊙ We¹ can think of the R matrix as the object that describes a pair of particles with two momenta.
- ⊙ If we say that the i^{th} particle lives in \mathcal{H}_i then the system lives in $\mathcal{H}_{\text{sys}} := \bigotimes_{k=1}^N \mathcal{H}_i$
- ⊙ The R matrix then is defined on two copies of \mathcal{H} - the state space of the two particle it describes.

 **Remark:** Often when describing the R matrix we do it in reference to two of three particles at a time e.g. R_{12} , R_{13} etc. The first one is easily represented by[1]

$$R_{12}(\lambda, \mu) = \underbrace{R(\lambda, \mu)}_{\mathcal{H}_1 \otimes \mathcal{H}_2} \otimes \underline{\mathbb{I}}_{\mathcal{H}_3} .$$

But the third and can't be written so easily - people write it in index form instead.

¹As physicists.

Coordinate Bethe Ansatz

The most intuitive form[1] of this is called the *coordinate* Bethe Ansatz - let's understand it.

$$\Psi(\{q\}) = \sum_{\sigma, \tau} \mathcal{A}(\sigma | \tau) \exp \left\{ i \sum_{j=1}^N q_{\sigma(j)} p_{\tau(j)} \right\} \Theta \left(q_{\sigma(1)} < q_{\sigma(2)} < \dots < q_{\sigma(N-1)} < q_{\sigma(N)} \right)$$